

Diffusion dynamics of QCD matter in nuclear collisions

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Hadron in Nucleus 2025

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Introduction

■ Exploring the QCD Phase diagram

QCD has a rich phase structure depending on the temperature and chemical potentials

Quark-gluon plasma
phase

(Critical point)

Hadronic phase

(Color superconductor)

Introduction

■ Where is the quark gluon plasma (QGP)?

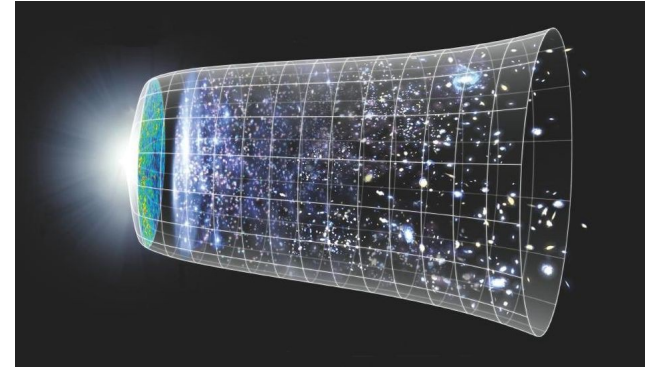
▶ Early universe

QGP had filled the early universe 10^{-5} - 10^{-4} s after the big bang

▶ Collider experiments ✓

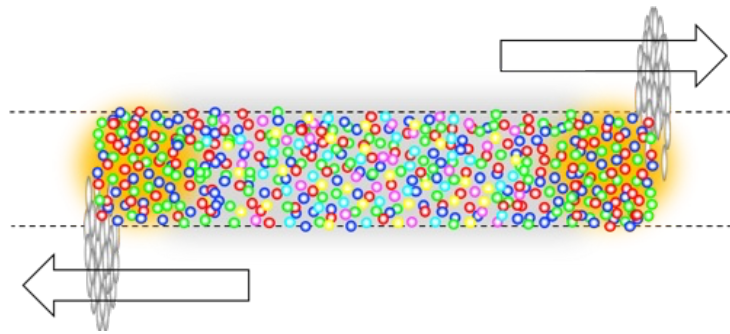
QGP can be produced by the collision of nuclei (especially heavy ions such as Au, Pb)

BNL RHIC, CERN LHC (SPS), etc.



Introduction

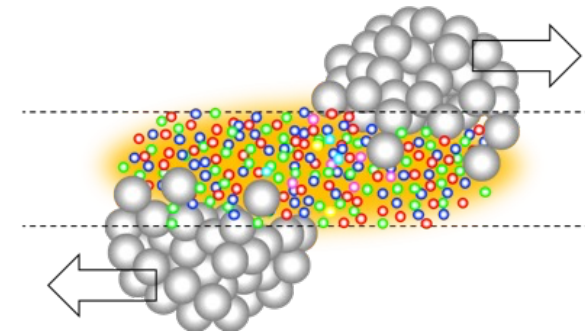
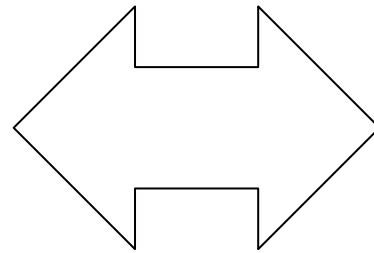
■ Exploring QCD matter at finite density



High-energy nuclear collisions
 $\sqrt{s_{NN}} \sim O(10^2 - 10^3) \text{ GeV}$

Dominated by particles from pair creation

→ High T , low μ_B



Low-energy nuclear collisions
 $\sqrt{s_{NN}} \sim O(1 - 10) \text{ GeV}$

Baryon density of nuclei is relevant

→ Low T , high μ_B

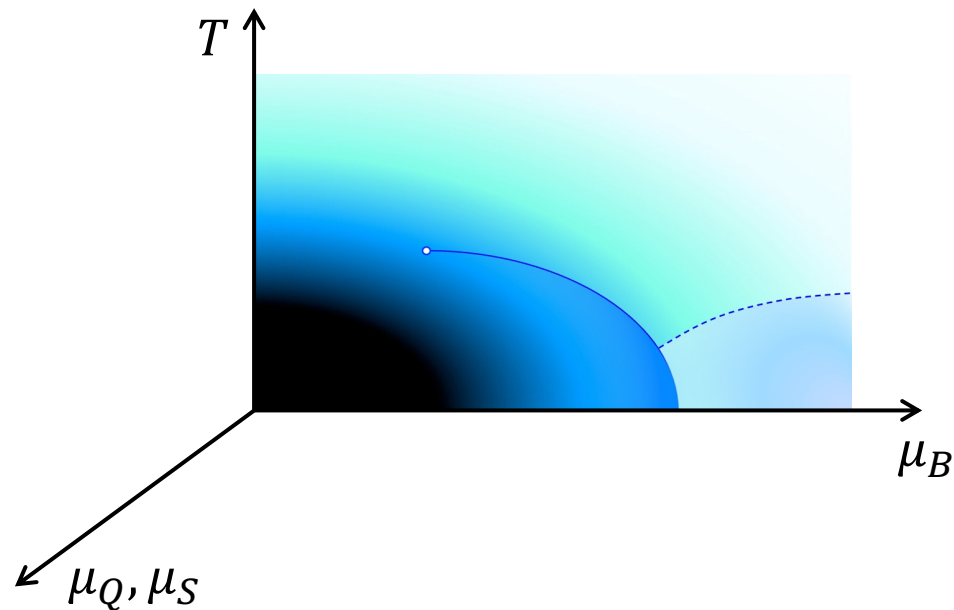
► Performed at RHIC-BES (& SPS), planned at FAIR-CBM, HIAF-CEE, J-PARC-HI, ...

Nuclear collisions

■ Conserved charges

The QGP in nuclear collisions are made of light quarks (u, d, s) ($T \sim 200$ MeV)

Baryon (B) **Electric charge (Q)** **Strangeness (S)** are conserved



The QCD phase diagram has to be extended to 4 dimensions

T : Temperature

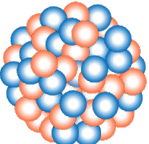
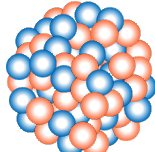
μ_B : Baryon chemical potential

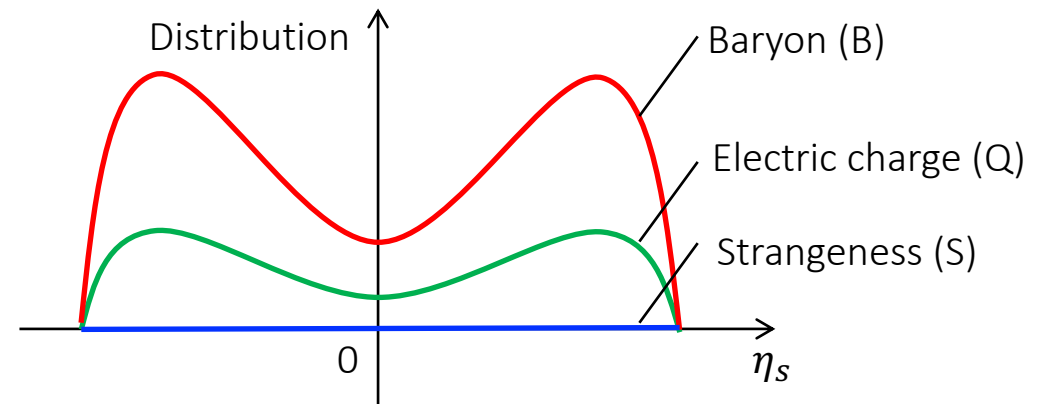
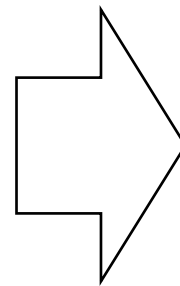
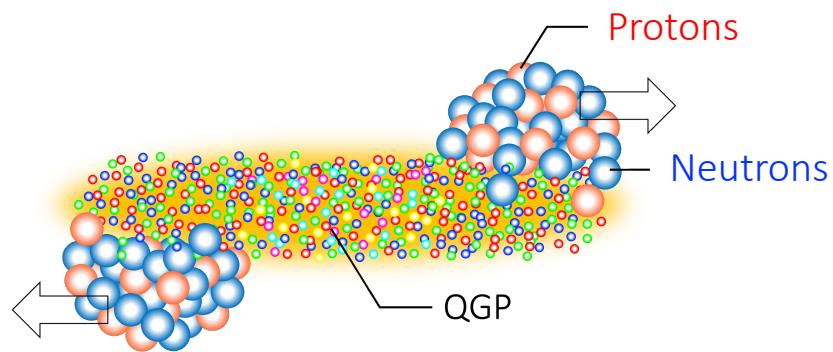
μ_Q : Charge chemical potential

μ_S : Strangeness chemical potential

Nuclear collisions

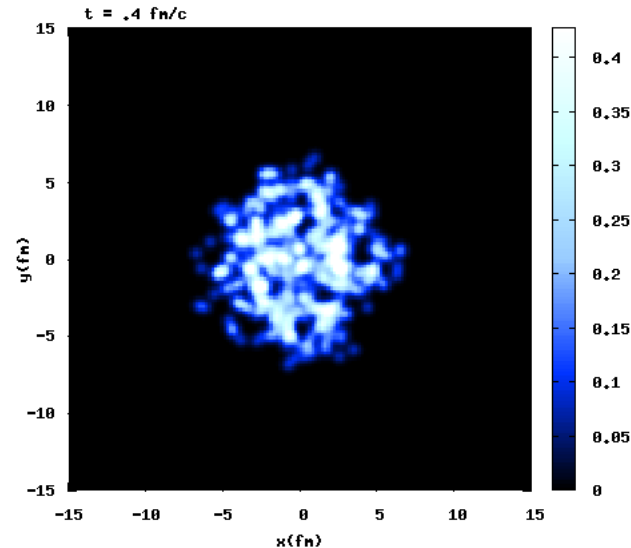
■ Geometry

^{197}Au	^{208}Pb
	
$Z = 79$ $N = 118$	$Z = 82$ $N = 126$
$\Rightarrow \frac{Q}{B} \approx 0.401$ $S = 0$	$\Rightarrow \frac{Q}{B} \approx 0.394$ $S = 0$

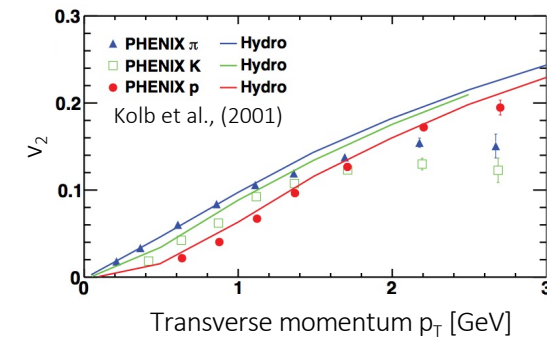
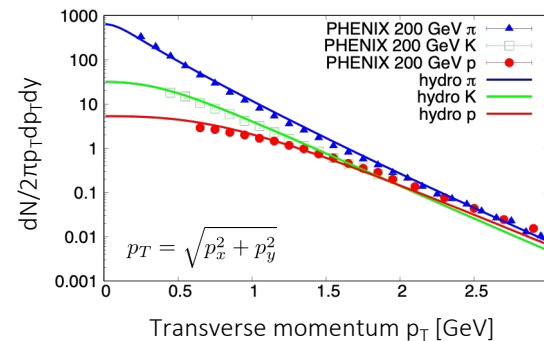


Nuclear collisions

■ Fluidity of the QCD matter



Relativistic hydrodynamic model describes nuclear collisions at high energies



We investigate diffusion dynamics of multiple conserved charges (B,Q,S) in the QCD matter

AM, Phys. Rev. C 86, 014908 (2012), Jan A. Fotakis et al., Phys. Rev. D 101, 076007 (2020)

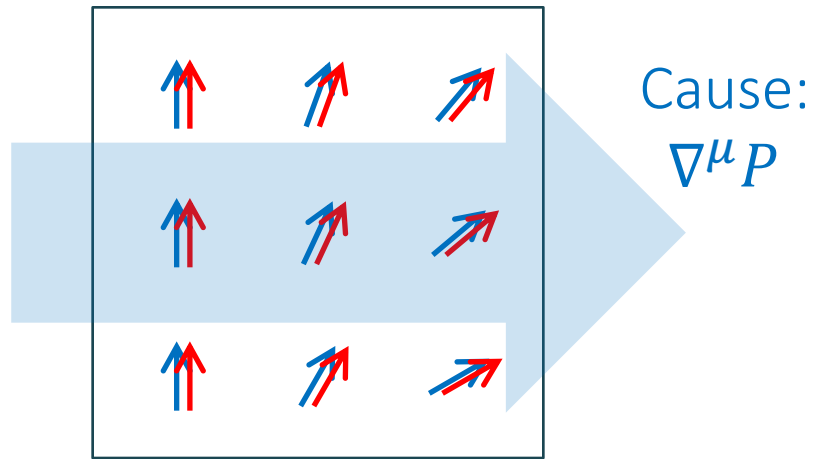


Hydrodynamic model

■ Difference between diffusion and convection

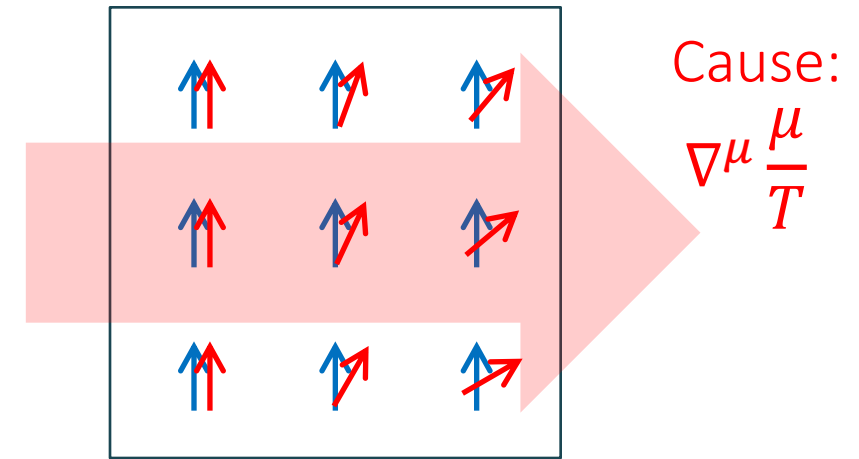
↑ flow u^μ
↑ conserved current N^μ

Convection



Charges are carried **along with** the flow
It can occur in ideal hydrodynamics
(**no** entropy production)

Diffusion



Charges are carried **against** the flow
Only in dissipative hydrodynamics
(**finite** entropy production)

Hydrodynamic model

■ Conserved currents

Energy-momentum tensor

$$T^{\mu\nu} = (e + P)u^\mu u^\nu - P g^{\mu\nu}$$

Net baryon current

$$N_B^\mu = n_B u^\mu + V_B^\mu$$

Net charge current

$$N_Q^\mu = n_Q u^\mu + V_Q^\mu$$

Net strangeness current

$$N_S^\mu = n_S u^\mu + V_S^\mu$$

*Shear and bulk viscous effects are orthogonal and omitted for simplicity

■ Equations of motion

$$\partial_\mu T^{\mu\nu} = 0, \quad \partial_\mu N_B^\mu = 0, \quad \partial_\mu N_Q^\mu = 0, \quad \partial_\mu N_S^\mu = 0$$

+ Constitutive relations

Hydrodynamic model

■ Relativistic diffusion equations

$$\begin{pmatrix} V_B^\mu \\ V_Q^\mu \\ V_S^\mu \end{pmatrix} = \underbrace{\mathcal{K} \nabla^\mu \begin{pmatrix} \mu_B/T \\ \mu_Q/T \\ \mu_S/T \end{pmatrix}}_{\text{Spatial derivatives of fugacities}} - \underbrace{\mathcal{T} D \begin{pmatrix} V_B^\mu \\ V_Q^\mu \\ V_S^\mu \end{pmatrix}}_{\text{Time derivative of diffusion currents}}$$

Conductivity: $\mathcal{K} = \begin{pmatrix} \kappa_{BB} & \kappa_{BQ} & \kappa_{BS} \\ \kappa_{QB} & \kappa_{QQ} & \kappa_{QS} \\ \kappa_{SB} & \kappa_{SQ} & \kappa_{SS} \end{pmatrix}$

Symmetric for Onsager reciprocal relations
Positive semi-definite for the 2nd law of thermodynamics

Relaxation time: $\mathcal{T} = \begin{pmatrix} \tau_{BB} & \tau_{BQ} & \tau_{BS} \\ \tau_{QB} & \tau_{QQ} & \tau_{QS} \\ \tau_{SB} & \tau_{SQ} & \tau_{SS} \end{pmatrix} = \tau_R I$

Assumed to be proportional to identity matrix for simplicity

Numerical setup

■ 1+1D hydrodynamic model

Initial conditions: ~ 17.3 GeV Pb+Pb collisions at 0-5% centrality

NA49, PRL 82, 2471 (1999)

$$\begin{aligned} e(\tau_{\text{hyd}}, \eta_s) &= a_1 \exp(-a_2 \eta_s^2 - a_3 \eta_s^4) \\ n_B(\tau_{\text{hyd}}, \eta_s) &= n_B^+(\eta_s) + n_B^-(\eta_s) \\ n_Q(\tau_{\text{hyd}}, \eta_s) &= 0.4 n_B(\tau_{\text{hyd}}, \eta_s) \\ n_S(\tau_{\text{hyd}}, \eta_s) &= 0 \end{aligned} \quad n_B^\pm(\eta_s) = \begin{cases} b_1 \exp[-b_2(\eta_s \mp \eta_0)^2 - b_3(\eta_s \mp \eta_0)^4] & (\eta_s > \pm \eta_0) \\ b_1 \exp[-\tilde{b}_2(\eta_s \mp \eta_0)^2 - \tilde{b}_3(\eta_s \mp \eta_0)^4] & (\eta_s \leq \pm \eta_0) \end{cases}$$

Hydrodynamization time: $\tau_{\text{hyd}} = 3 \text{ fm}/c$

Transport coefficients: $\frac{\kappa_B}{T^2} = 0.05$, $\frac{\kappa_Q}{T^2} = 0.02$, $\frac{\kappa_S}{T^2} = 0.1$

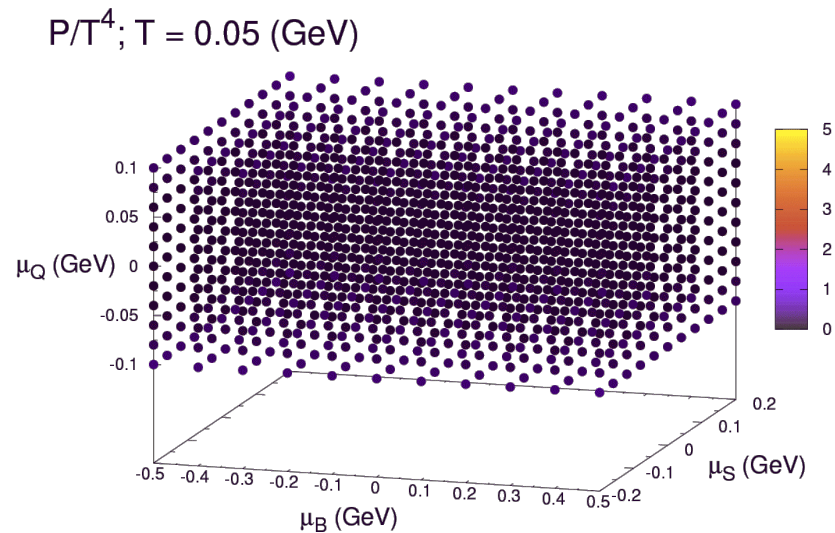
$$\kappa_{BQ} = c\sqrt{\kappa_B \kappa_Q}, \quad \kappa_{BS} = c\sqrt{\kappa_B \kappa_S}, \quad \kappa_{QS} = c\sqrt{\kappa_Q \kappa_S} \quad (c = 0 \text{ or } 0.5)$$

Numerical setup

■ 1+1D hydrodynamic model

Equation of state: Lattice QCD + hadron resonance gas model (NEOS-4D)

AM, G. Pihan, B. Schenke, C. Shen, Phys. Rev. C 110, 044905 (2024)



A full 4-dimensional EoS with B, Q, S charges in the range relevant to nuclear collisions



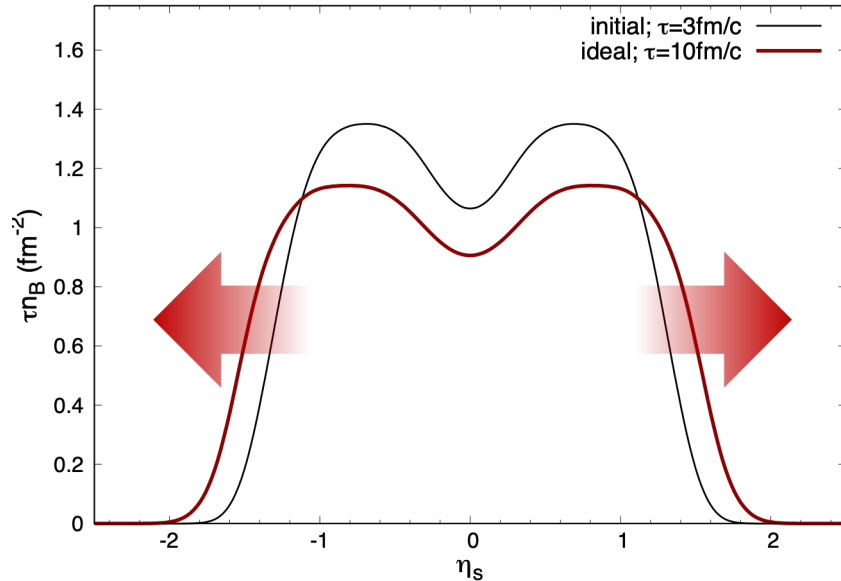
The EoS is available for downloading:

<https://sites.google.com/view/qcdneos4d/home>



Results

■ Spacetime rapidity distribution of net baryon density (n_B)

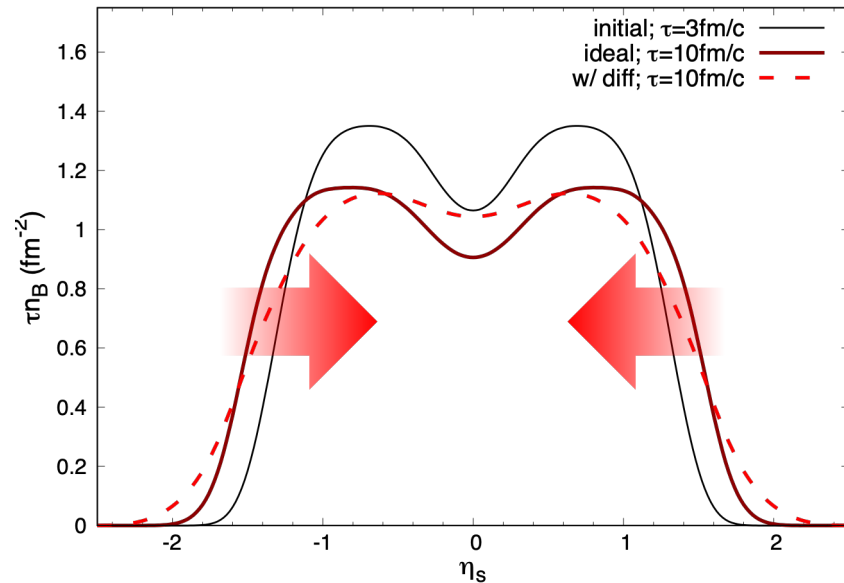


1. Ideal hydrodynamic evolution

Carried to forward rapidity by convection

Results

■ Spacetime rapidity distribution of net baryon density (n_B)



1. Ideal hydrodynamic evolution

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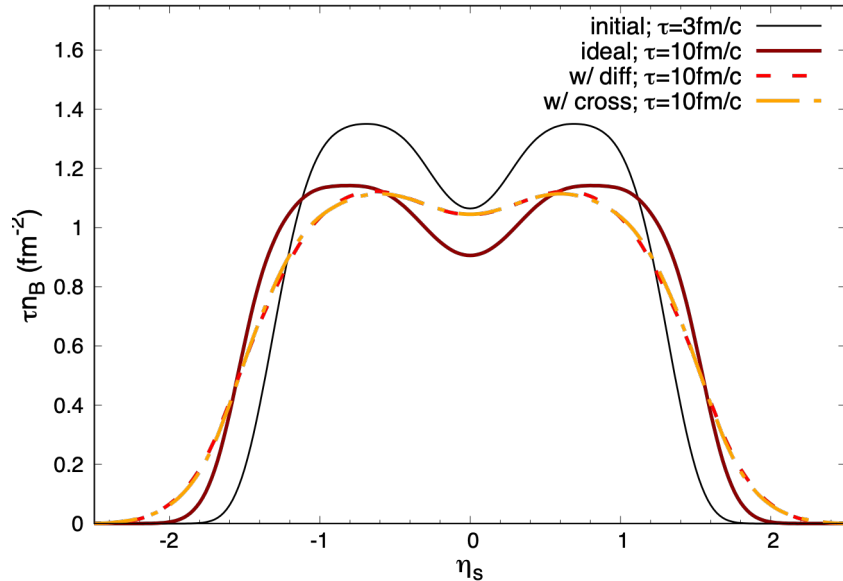
2. With diffusion currents

Partially brought back to mid-rapidity by gradients in μ_B/T

$$V_B^\mu \approx \kappa_{BB} \nabla^\mu \frac{\mu_B}{T} + \kappa_{BQ} \nabla^\mu \frac{\mu_Q}{T} + \kappa_{BS} \nabla^\mu \frac{\mu_S}{T}$$

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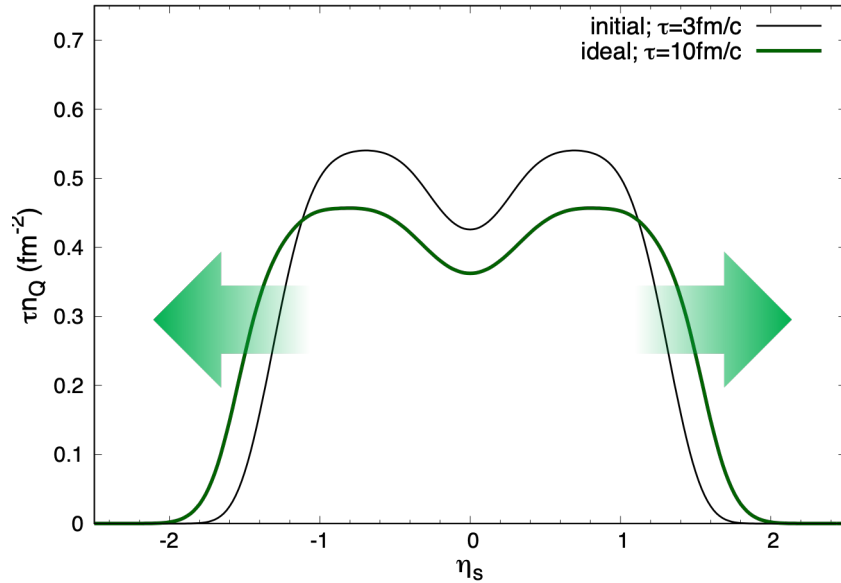
Partially brought back to mid-rapidity by gradients in μ_B/T

3. With cross-coupling currents

Effects of the other charges are negligible

Results

■ Spacetime rapidity distribution of net charge density (n_Q)

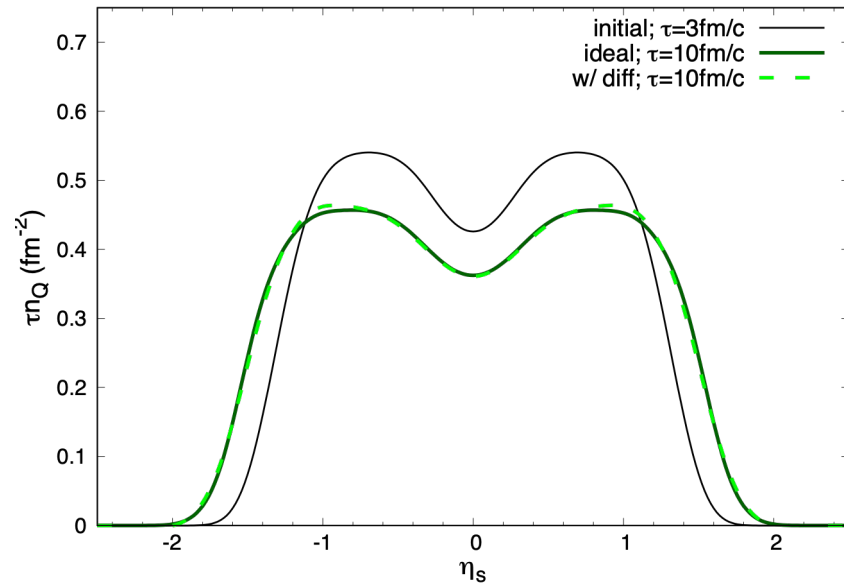


1. Ideal hydrodynamic evolution

Carried to forward rapidity by convection

Results

■ Spacetime rapidity distribution of net charge density (n_Q)



1. Ideal hydrodynamic evolution

Carried to forward rapidity by convection

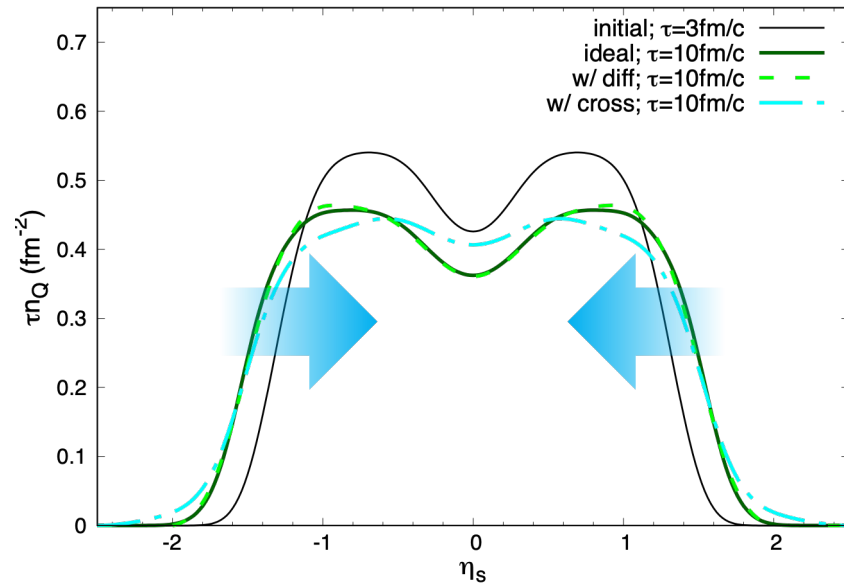
2. With diffusion currents

Carried to forward rapidity, but the effect is small because μ_Q is small

$$V_Q^\mu \approx \kappa_{BQ} \nabla^\mu \frac{\mu_B}{T} + \kappa_{QQ} \nabla^\mu \frac{\mu_Q}{T} + \kappa_{QS} \nabla^\mu \frac{\mu_S}{T}$$

Results

■ Spacetime rapidity distribution of net charge density (n_Q)



$$V_Q^\mu \approx \kappa_{BQ} \nabla^\mu \frac{\mu_B}{T} + \kappa_{QQ} \nabla^\mu \frac{\mu_Q}{T} + \kappa_{QS} \nabla^\mu \frac{\mu_S}{T}$$

1. Ideal hydrodynamic evolution

Carried to forward rapidity by convection

2. With diffusion currents

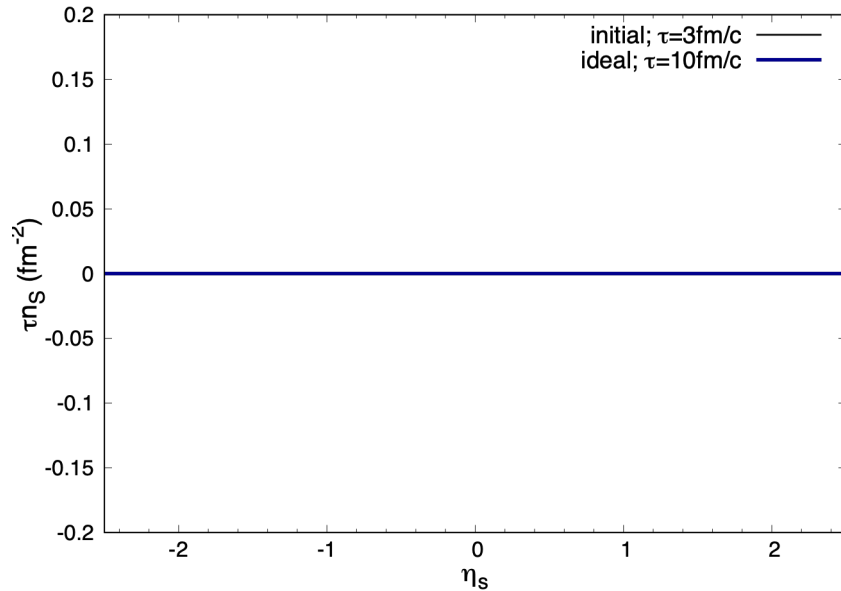
Carried to forward rapidity, but the effect is small because μ_Q is small

3. With cross-coupling currents

Carried back to mid-rapidity by μ_B

Results

■ Spacetime rapidity distribution of net strangeness density (n_s)

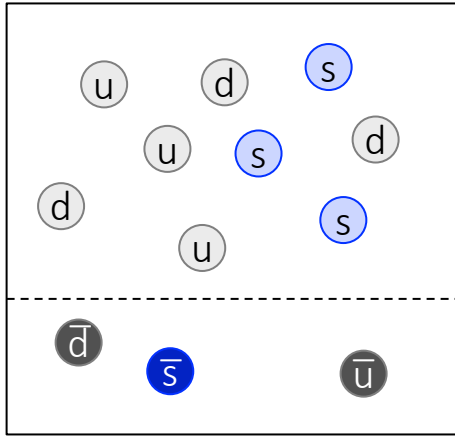


1. Ideal hydrodynamic evolution

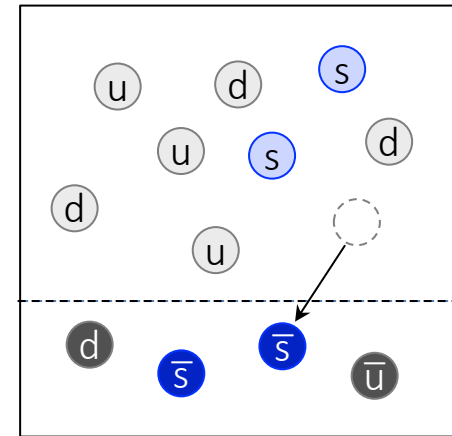
Zero everywhere during time evolution

Discussion

■ Strangeness neutrality condition



$\mu_B > 0$ and $\mu_S = 0$
means $n_S \neq 0$

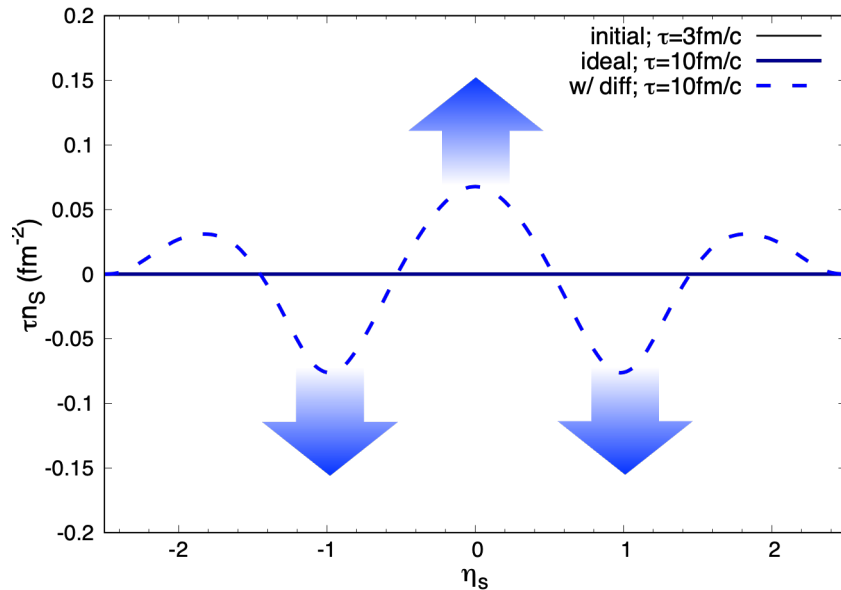


$\mu_B > 0$ and $n_S = 0$
leads to $\mu_S > 0$

s quark chemical potential $\mu_s = \frac{1}{3}\mu_B - \frac{1}{3}\mu_Q - \mu_S = 0$ implies $\mu_S > 0$
when $\mu_B > 0$ and $\mu_Q < 0$

Results

■ Spacetime rapidity distribution of net strangeness density (n_s)



1. Ideal hydrodynamic evolution

Zero everywhere during time evolution

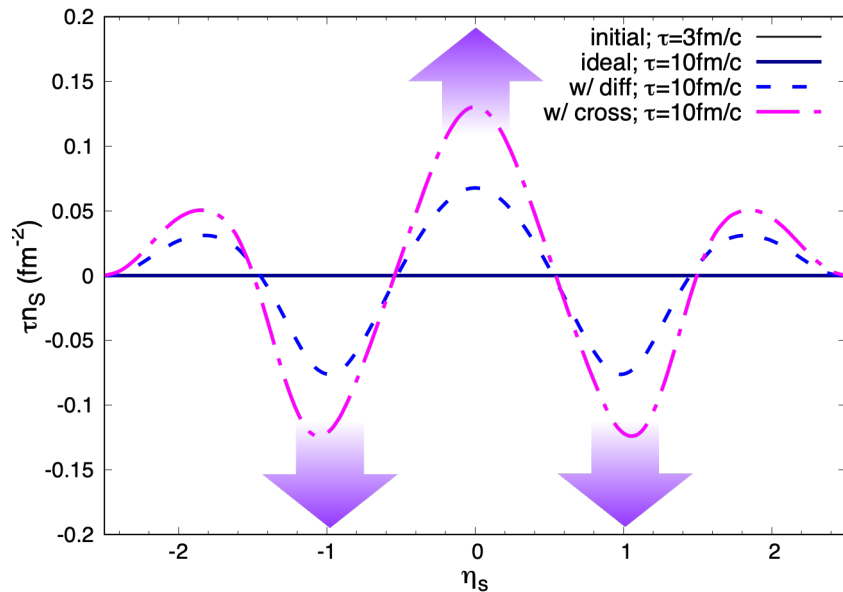
2. With diffusion currents

Gradients in μ_s/T induce non-zero n_s
(integrated strangeness is zero)

$$V_S^\mu \approx \kappa_{BS} \nabla^\mu \frac{\mu_B}{T} + \kappa_{QS} \nabla^\mu \frac{\mu_Q}{T} + \kappa_{SS} \nabla^\mu \frac{\mu_S}{T}$$

Results

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Zero everywhere during time evolution

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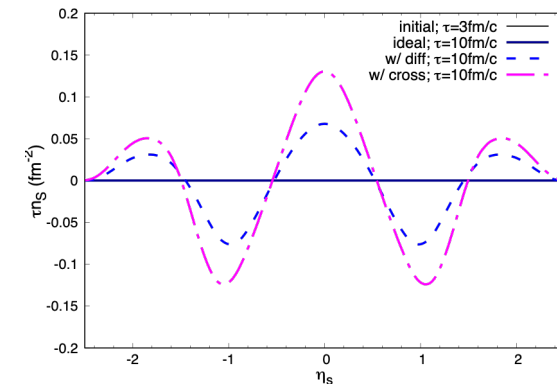
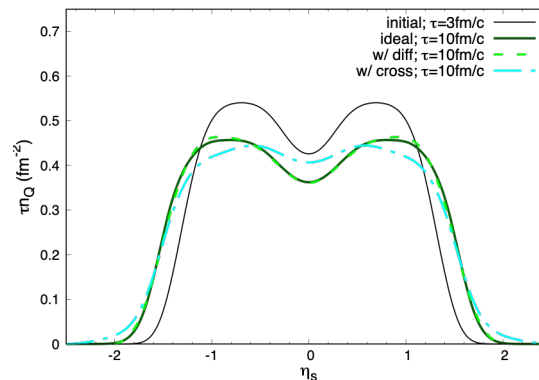
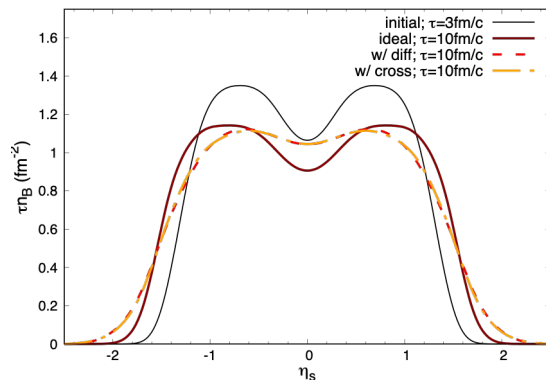
Gradients in μ_S/T induce non-zero n_s
(integrated strangeness is zero)

3. With cross-coupling currents

The amplitude is modified by μ_B

Summary and outlook

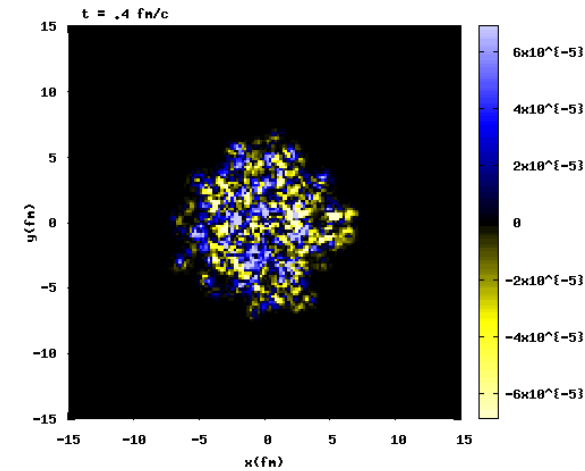
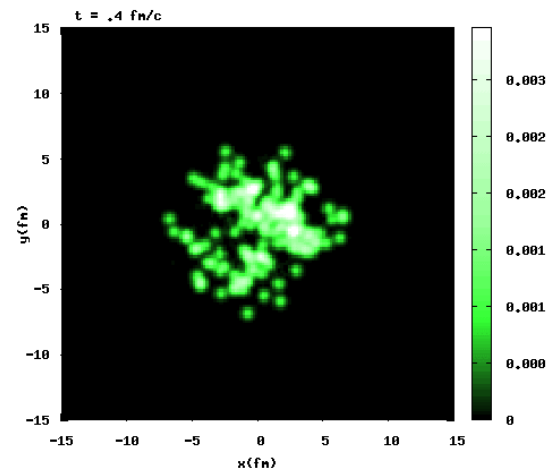
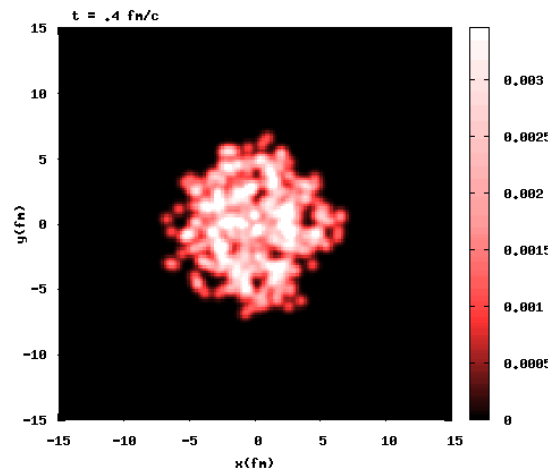
- Diffusion currents of **net baryon**, **electric charge**, and **strangeness** are introduced in the hydrodynamic model of nuclear collisions
- ▶ **Baryon diffusion**: n_B is carried inwards; cross-coupling effect is minute
 - ▶ **Charge diffusion**: Small effect; cross-coupling effect is large
 - ▶ **Strangeness diffusion**: Non-zero n_S is induced (the neutrality condition is broken); cross-coupling effect is large



Summary and outlook

■ Prospects

- ▶ Full 3D simulations by introducing transverse dynamics to estimate particle spectra and elliptic flow
- ▶ Estimations of initial conditions and transport coefficients, etc.



The end

- Thank you for listening!